

# Dark sector physics search in missing energy events with the NA64 experiment

B. Radics

(on behalf of the NA64 collaboration)



# Outline

- Motivation and method of search
- The NA64 experiment
- Runs in 2016 and 2017
- Simulation of Dark Photon production
- Analysis of data
- Results
- Conclusions

# Motivation

- Possible candidates for new physics: sub-GeV dark sector particles not charged under SM forces, only gravitational interaction, "portal" interactions with SM particles
- Thermal freeze-out of DM-SM could explain relic density, and put constraints on the parameter space
- May affect galactic structure formation, muon  $(g-2)_\mu$ , etc
- Parameter space is poorly tested
- Most accessible via portal interactions with SM: gauge kinetic mixing, MeV - GeV mass range, high intensity searches
- Most viable is interaction of DM with SM through a vector portal  $A'$  boson

Dark Sectors 2016 Workshop: Community Report, J.Alexander et al., arxiv: 1608.8632

# Motivation

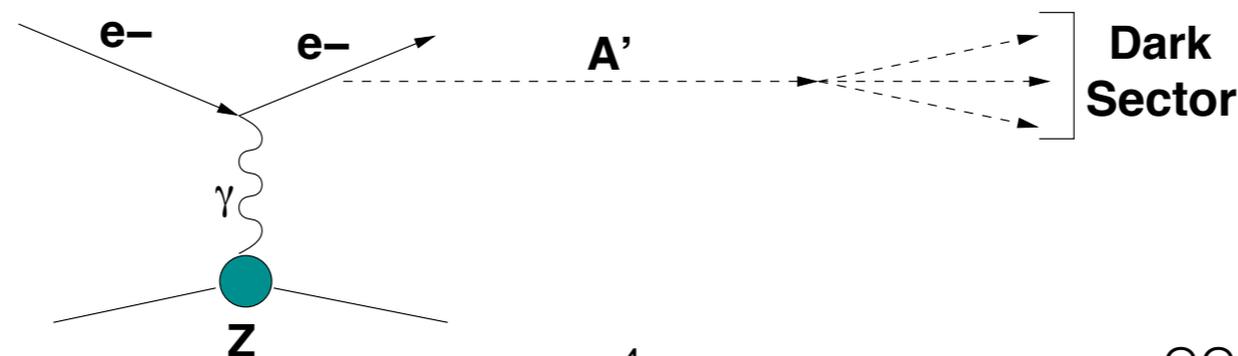
- New  $A'$  vector portal boson (dark photon) could mix kinetically with photon

$$\mathcal{L} = \mathcal{L}_{\text{SM}} - \frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + \frac{\epsilon}{2} F'_{\mu\nu} F^{\mu\nu} + \frac{m_{A'}^2}{2} A'_\mu A'^\mu + i\bar{\chi}\gamma^\mu\partial_\mu\chi - m_\chi\bar{\chi}\chi - e_D\bar{\chi}\gamma^\mu A'_\mu\chi$$

- $A'$  corresponds to new  $U(1)_D$  gauge symmetry,  $\epsilon \ll 1$
- Requirement of thermal freeze-out of DM-SM annihilation through photon- $A'$  mixing allows to derive relations between the parameters (PRD 91,094026 (2015)).
- Rate of DM annihilation into SM fermions, allows to define signal event rate,  $y$ ,

$$\langle\sigma v\rangle \propto \underbrace{\alpha_{\text{DM}}\epsilon^2 (m_\chi^4/m_{A'}^4)}_y \alpha/m_\chi^2 \quad \alpha_{\text{DM}} = e_D^2/4\pi$$

- Decay channels: visible:  $e^+e^-$ ,  $\mu^+\mu^-$ , hadron, ..., invisible:  $A' \rightarrow \chi\chi^-$  if  $m_{A'} > 2m_\chi$ . It is dominante if  $\alpha_{\text{DM}} \gg \epsilon$ .
- Production: interaction of high energy electrons in an active beam dump target



# NA64 collaboration

D. Banerjee,<sup>11</sup> V. Burtsev,<sup>9</sup> D. Cooke,<sup>11</sup> P. Crivelli,<sup>11</sup> E. Depero,<sup>11</sup> A. V. Dermenev,<sup>4</sup> S. V. Donskov,<sup>8</sup> F. Dubinin,<sup>5</sup>  
R. R. Dusaev,<sup>9</sup> S. Emmenegger,<sup>11</sup> A. Fabich,<sup>3</sup> V. N. Frolov,<sup>2</sup> A. Gardikiotis,<sup>7</sup> S. N. Gninenko\*,<sup>4</sup> M. Hösgen,<sup>1</sup>  
V. A. Kachanov,<sup>8</sup> A. E. Karneycu,<sup>4</sup> B. Ketzner,<sup>1</sup> D. V. Kirpichnikov,<sup>4</sup> M. M. Kirsanov,<sup>4</sup> I. V. Konorov,<sup>5</sup>  
S. G. Kovalenko,<sup>10</sup> V. A. Kramarenko,<sup>6</sup> L. V. Kravchuk,<sup>4</sup> N. V. Krasnikov,<sup>4</sup> S. V. Kuleshov,<sup>10</sup> V. E. Lyubovitskij,<sup>9</sup>  
V. Lysan,<sup>2</sup> V. A. Matveev,<sup>2</sup> Yu. V. Mikhailov,<sup>8</sup> V. V. Myalkovskiy,<sup>2</sup> V. D. Peshekhonov<sup>†,2</sup> D. V. Peshekhonov,<sup>2</sup>  
O. Petuhov,<sup>4</sup> V. A. Polyakov,<sup>8</sup> B. Radics,<sup>11</sup> A. Rubbia,<sup>11</sup> V. D. Samoylenko,<sup>8</sup> V. O. Tikhomirov,<sup>5</sup> D. A. Tlisov,<sup>4</sup>  
A. N. Toropin,<sup>4</sup> A. Yu. Trifonov,<sup>9</sup> B. Vasilishin,<sup>9</sup> G. Vasquez Arenas,<sup>10</sup> P. Ulloa,<sup>10</sup> K. Zhukov,<sup>5</sup> and K. Zioutas<sup>7</sup>  
(The NA64 Collaboration<sup>†</sup>)

<sup>1</sup>*Universität Bonn, Helmholtz-Institut für Strahlen-und Kernphysik, 53115 Bonn, Germany*

<sup>2</sup>*Joint Institute for Nuclear Research, 141980 Dubna, Russia*

<sup>3</sup>*CERN, European Organization for Nuclear Research, CH-1211 Geneva, Switzerland*

<sup>4</sup>*Institute for Nuclear Research, 117312 Moscow, Russia*

<sup>5</sup>*P.N. Lebedev Physics Institute, Moscow, Russia, 119 991 Moscow, Russia*

<sup>6</sup>*Skobel'syn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia*

<sup>7</sup>*Physics Department, University of Patras, Patras, Greece*

<sup>8</sup>*State Scientific Center of the Russian Federation Institute for High Energy Physics of National Research Center 'Kurchatov Institute' (IHEP), 142281 Protvino, Russia*

<sup>9</sup>*Tomsk Polytechnic University, 634050 Tomsk, Russia*

<sup>10</sup>*Universidad Técnica Federico Santa María, 2390123 Valparaíso, Chile*

<sup>11</sup>*ETH Zürich, Institute for Particle Physics, CH-8093 Zürich, Switzerland*

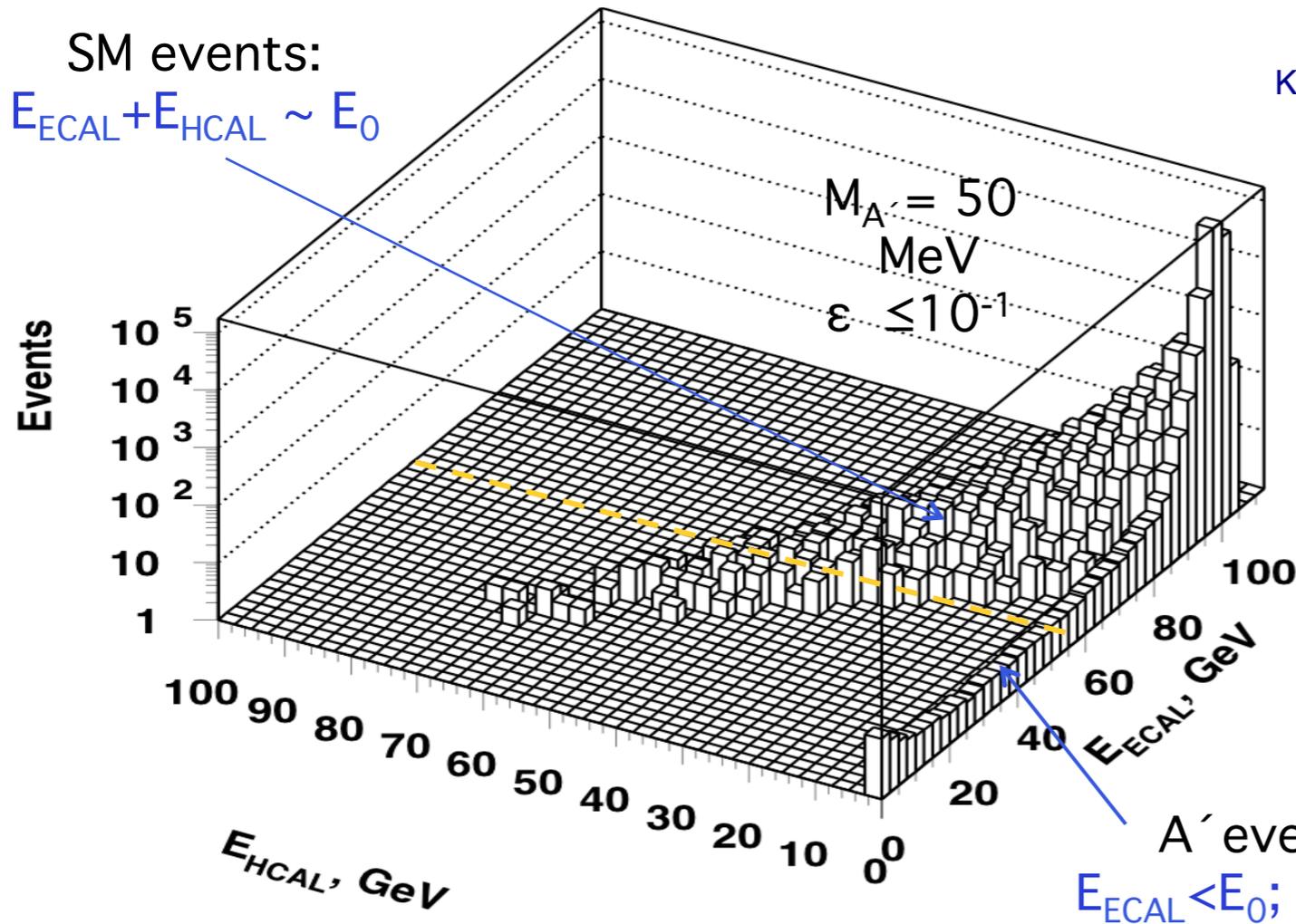
47 researchers from 11 institutes  
Proposed in 2014, first test beam in 2015

# Method of search for $A'$ $\rightarrow$ invisible

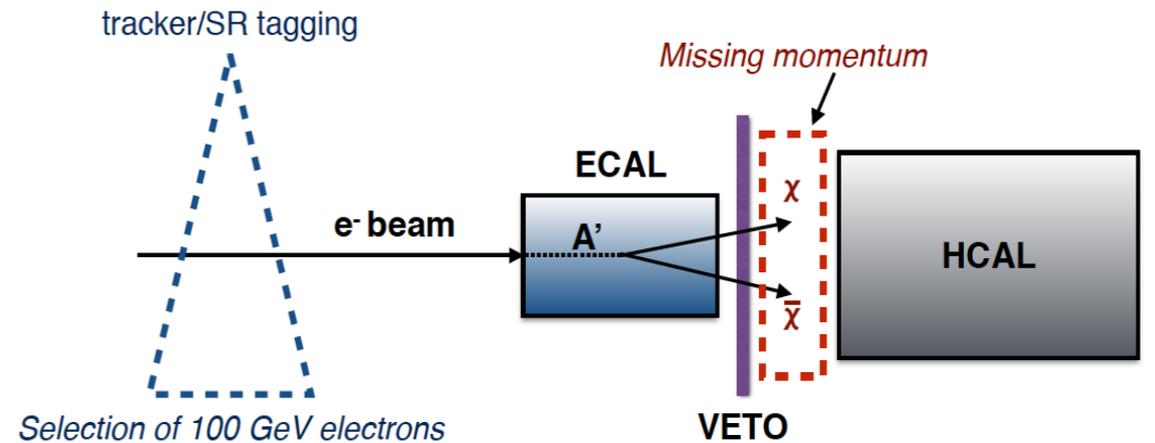
- If realised by nature, any source of photons will produce all kinematically possible massive  $A'$  states with the appropriate mixing strength: e.g. kinetic mixing with bremsstrahlung photons in the reaction of high-energy electrons from a beam absorbed in an active beam dump.
- Followed by the prompt decay  $A' \rightarrow$  invisible into DM particles:  $e^- Z \rightarrow e^- Z A'$ ;  $A' \rightarrow \chi \chi^-$
- A fraction of the beam energy,  $f$ , is carried away by  $\chi$  particles, penetrating the target without interactions,  $E_{A'} = f E_0$
- The remaining part of the beam energy is deposited in the target:  $E_e = (1-f) E_0$
- Signal signature: excess of events above background with
  - single isolated energy e-m shower with energy  $E_e < E_0$
  - missing energy  $E_{\text{miss}} = E_{A'} = E_0 - E_e$
- Number of  $A'$  produced per electron on target (EOT):

$$n_{A'}(\epsilon, m_{A'}, E_0) = \frac{\rho N_A}{A_{Pb}} \sum_i n(E_0, E_e, s) \sigma^{A'}(E_e) \Delta s_i$$

# Simulation of $eZ \rightarrow eZ A'$ ; $A' \rightarrow$ invisible



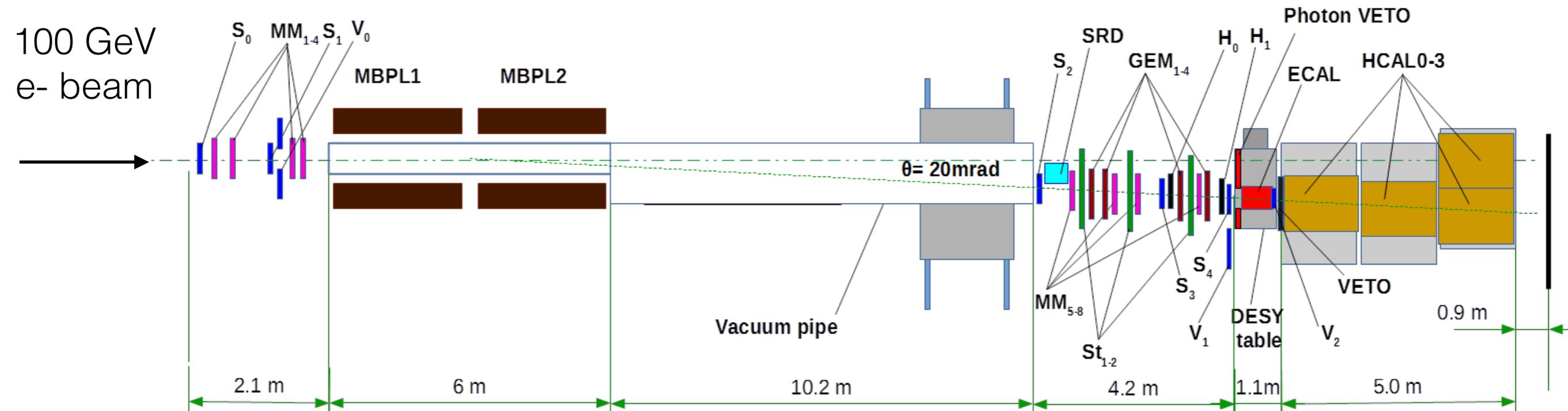
Gninenko, Kirsanov,  
 Krasnikov, Kirpichnikov  
 PRD(2016)



- Geant4 and  $A'$  emission in the e-m shower development.
- Cross section from Bjorken et al. 2009.
- Sensitivity  $\sim \epsilon^2$  ( $A'$  production vertex) - while for beam dump experiments  $\sim \epsilon^2 \alpha_D$  (+  $A'$  decay and  $\chi$  scattering off electrons in the target detector).
- For small  $\epsilon$  mixing parameter this scheme has great advantage.

# NA64 experiment setup invisible search mode

## TOP VIEW

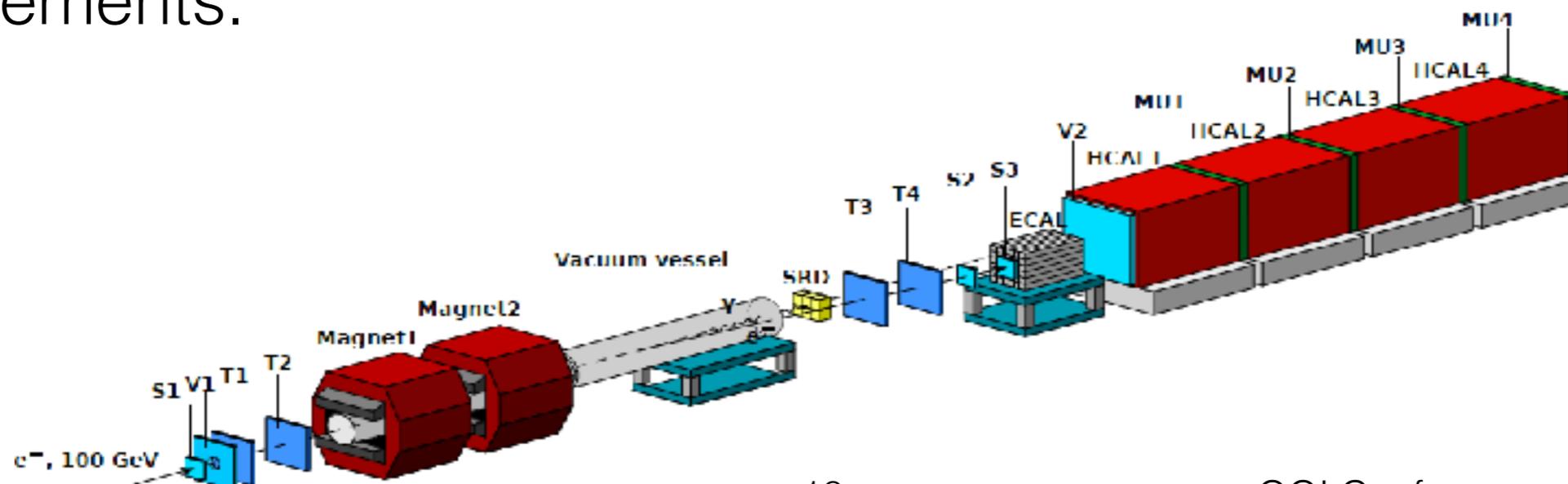


# NA64 experiment setup



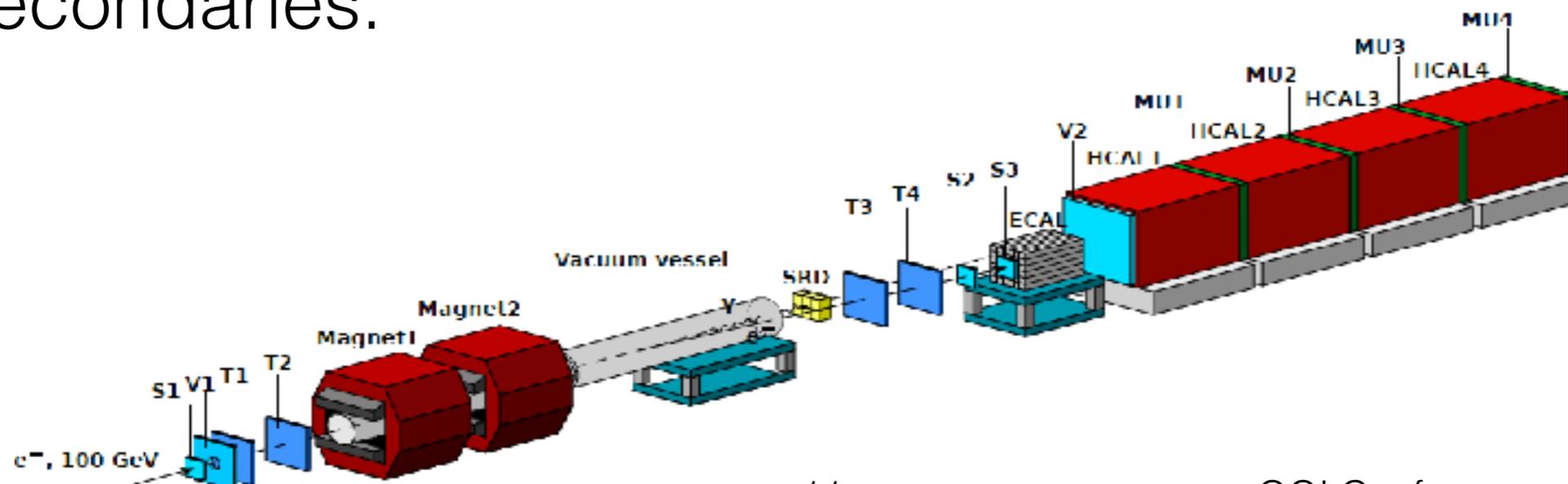
# Key moments in reconstruction

- Synchrotron Radiation detector (SRD) made as lead - scintillator sandwich used to suppress pions and other heavier than  $e^-$  particles from the beam.
- The shower profile in the ECAL is compared to profile of true electrons in order to suppress wrong particles.
- Micromegas track detectors are used to reconstruct the momentum of  $e^-$  before the ECAL to suppress small fraction of soft electrons from interaction in beam line elements.



# Key moments in reconstruction

- Each ECAL module is  $40 X_0$  with a  $4X_0$  preshower initial part, electron energy resolution:  $dE/E \sim 0.1/\sqrt{E}$
- Requiring in-time between SRDs combined with ECAL longitudinal and lateral shower information:  $\pi/e^- < 10^{-5}$ , 95%  $e^-$  ID efficiency (NIM A 866 (2017) 196).
- V2 after ECAL to veto charged secondaries, and HCAL ( $30 \lambda_{\text{int}}$ , Fe+Sc) to veto on muons or hadronic secondaries.

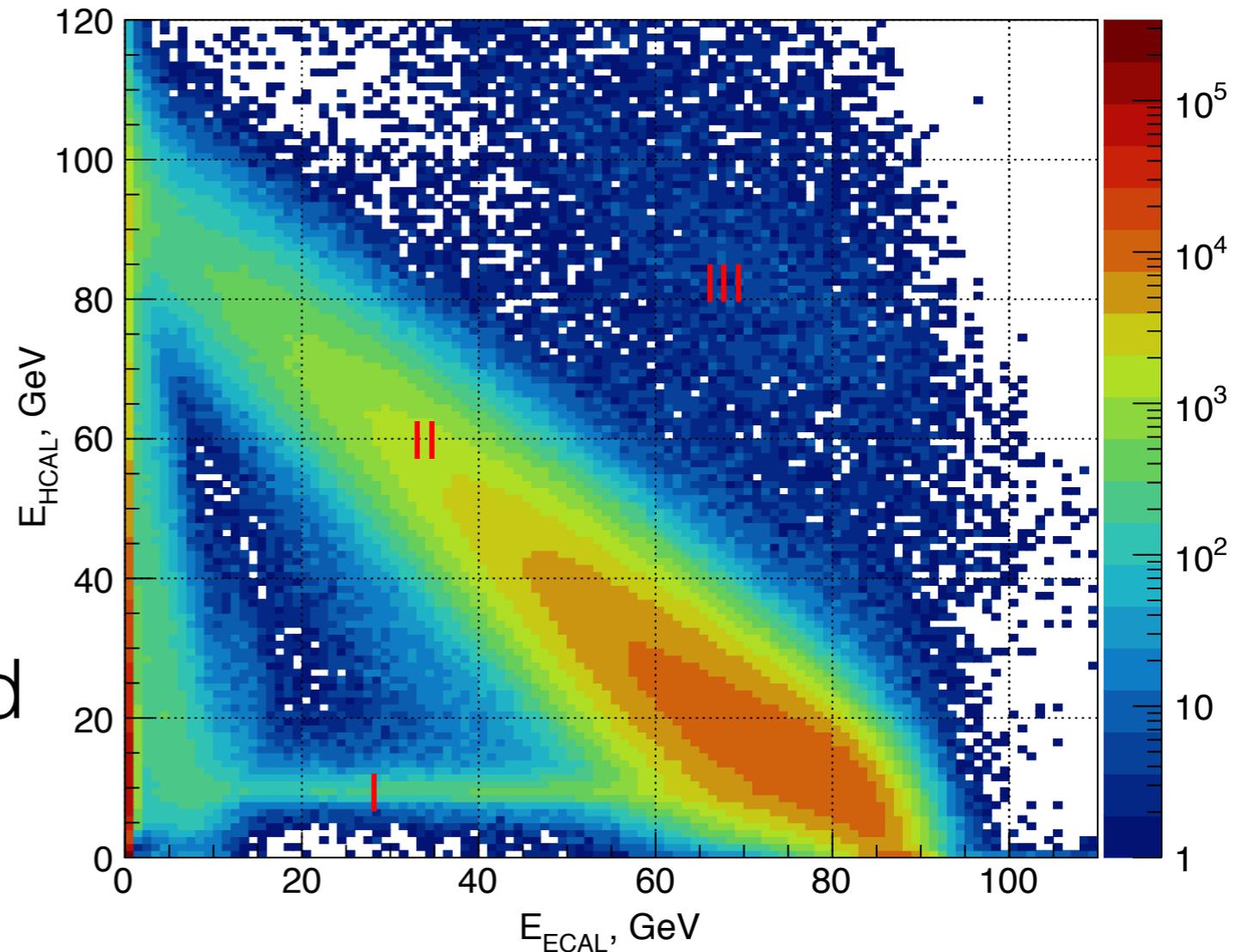


# Data taking in 2016

- 1st Run period: 29.06-13.07 (2w)
- 2nd Run period: 12.10-09.11 (4w)
  - Low intensity:  $n_{\text{EOT}} = 2.3 \times 10^{10}$  ( $\sim 1.4\text{-}2 \times 10^6$  e- /spill)
  - Medium intensity:  $n_{\text{EOT}} = 1.1 \times 10^{10}$  ( $\sim 3\text{-}3.5 \times 10^6$  e- /spill)
  - High intensity:  $n_{\text{EOT}} = 0.9 \times 10^{10}$  ( $\sim 4.5\text{-}5 \times 10^6$  e- /spill)
- $\text{Tr}(A') = \prod S_i \times V1 \times \text{PS}( > E_{\text{PS}} ) \times \text{ECAL}( < E_{\text{ECAL}} )$

# ECAL vs HCAL energy

- Region I: dimuon events
- Region II:  $E_{\text{ECAL}} + E_{\text{HCAL}} = 100 \text{ GeV}$
- Region III: pile-up of e- and beam hadrons (1-20%)

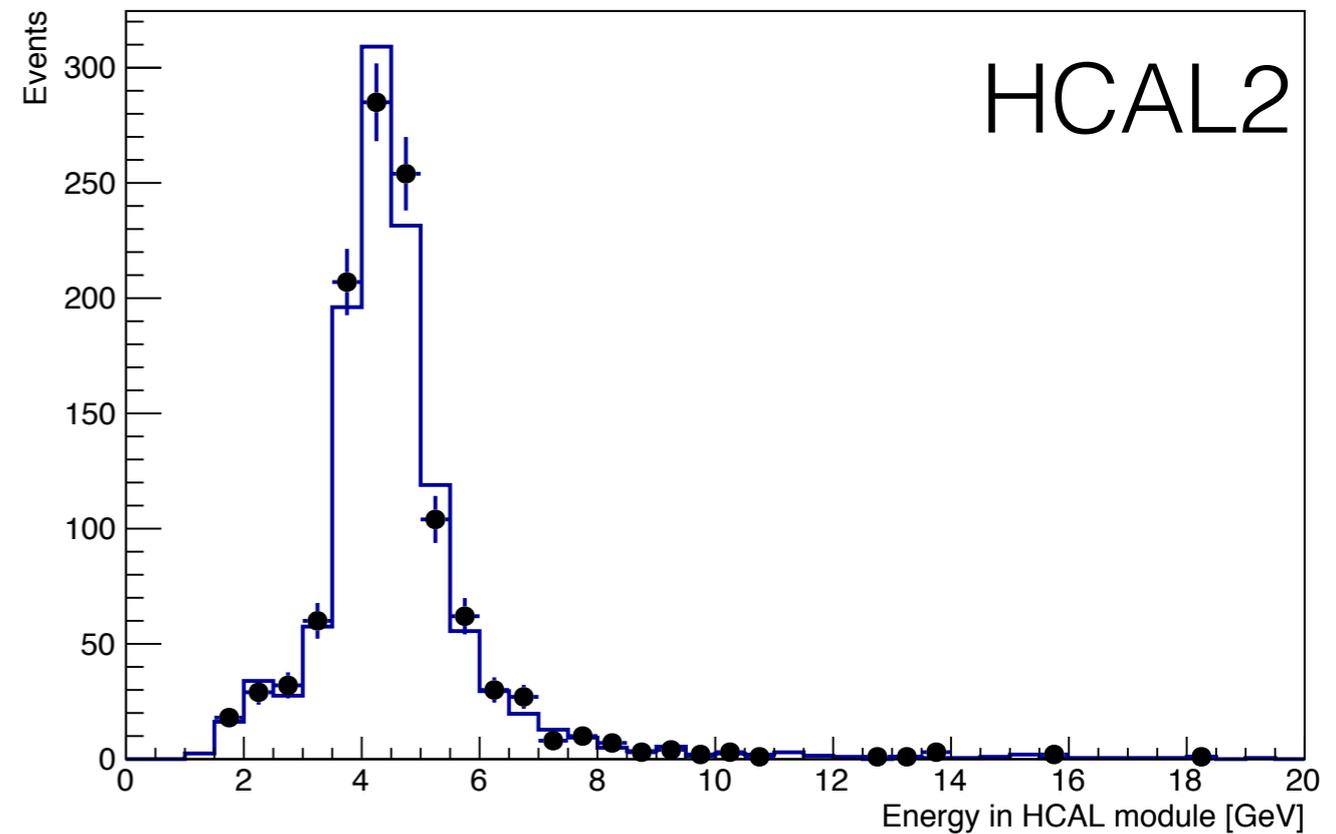
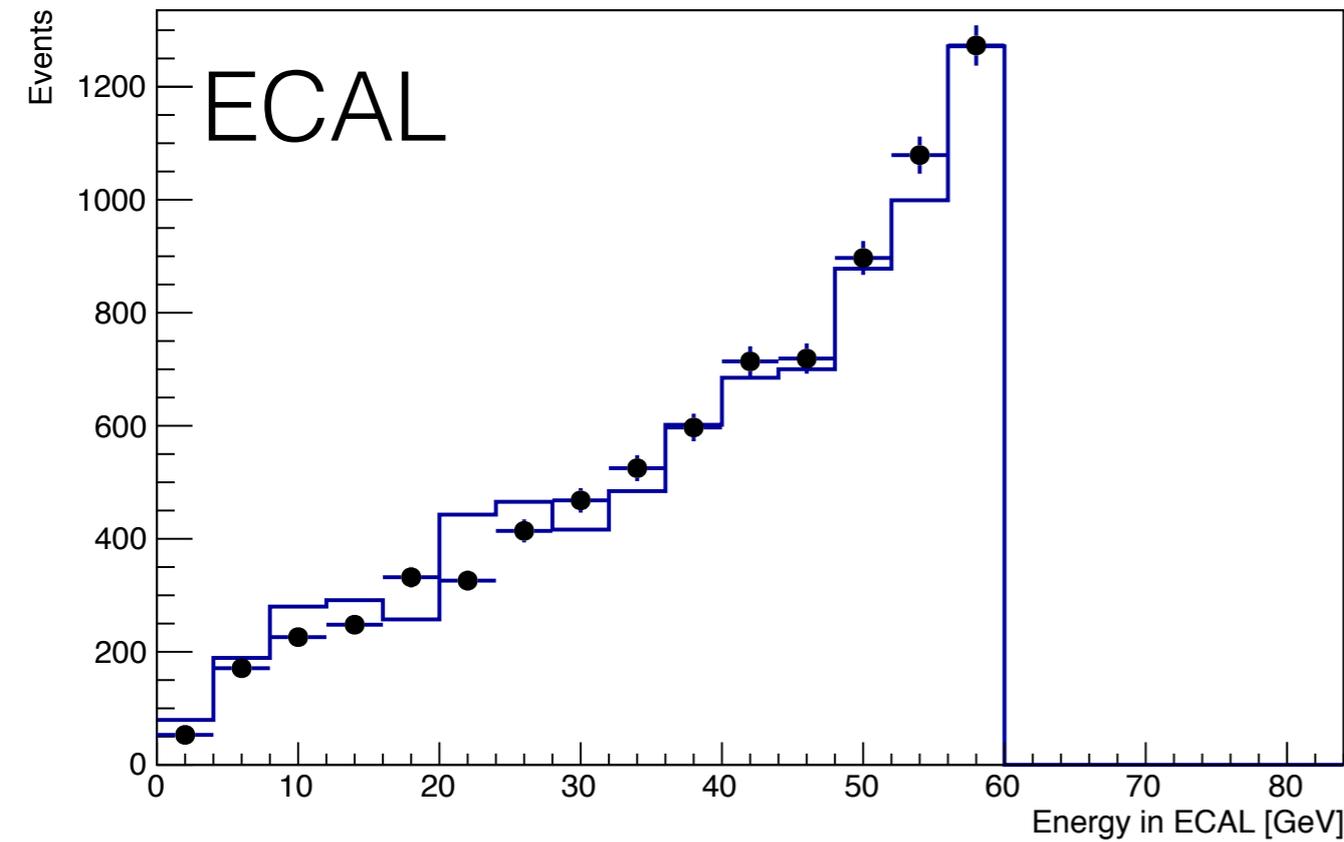


Only SRD selection to  
be e- event

# Dimuon production as reference

- Rare process gamma to muon conversion ( $eZ \rightarrow eZ\gamma; \gamma \rightarrow \mu\mu$ ), many similarities with our signal. Available in G4, off by default.
- Can be used to estimate corrections to signal reconstruction efficiency and uncertainties in  $A'$  yield calculations
- HCAL energy around 10 GeV.
- $\sim 10^4$  dimuon pairs detected in HCAL in 2016 run period.
- MC simulation: cross section have been biased in G4 by a factor of 200 to have good statistics.
- MC compared with Data.

# Dimuon reconstruction



Data sample	beam intensity, $10^6$	$n_{tot}, 10^6$	$n_{2\mu}^{MC}$	$n_{2\mu}^{data}$	Efficiency
run I	1.8	171	1223	1124	0.92
run II	3.2	208.5	1491	1268	0.85
run III	4.6	597	4271	3417	0.81

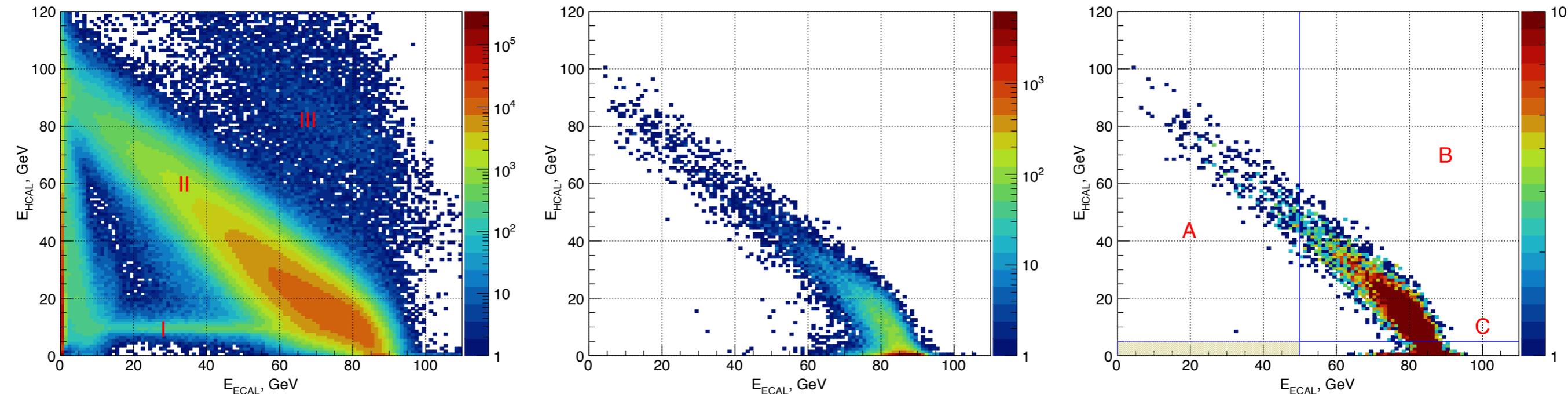
# Analysis: efficiency and uncertainties

Efficiency	Value, uncertainty	sample
number of collected EOT, $n_{EOT}$	$1 \pm 0.02$	$e^-$ Data
incoming $e^-$ selection cuts, $\epsilon_e$	$0.58 \pm 0.03$	$e^-$ Data
$A'$ yield, $\epsilon_{A'}$	$\epsilon, m_{A'}$ dependent, 10%	MC, Dimuons
ECAL selection cuts, $\epsilon_{ECAL}$	$0.93 \pm 0.06$	Data, Dimuons
Veto cut, $\epsilon_V$	$0.94 \pm 0.03$	Data, MC
HCAL selection cuts, $\epsilon_{HCAL}$	$0.98 \pm 0.02$	Data, MC
Total	$0.50 \pm 0.13$	

- Values correspond to high-intensity run.
- Total efficiency varying  $0.73 \pm 0.12$  to  $0.50 \pm 0.13$ .
- ECAL and incoming  $e^-$  selection most rate dependent.

# Analysis cuts

Medium beam intensity



- Left: only SRD cut to be e- events
- Middle: all selection but cut against upstream interactions (Tracker hit multiplicity, and lateral energy spread and time spread in HCAL cells)
- Right: final event selection

# Backgrounds

- Leak of energy through holes, cracks in the detector
  - X-Y scan of ECAL and HCAL - no significant E leak found
- Detector hermeticity: photo-nuclear reaction producing neutrons, charged hadrons escaping detection in HCAL (non-herm)
  - pion beam test, Data-MC comparison, single hadron prod. prob.  $< 10^{-4}$ , non hermeticity  $< 10^{-9}$ , overall negligible  $< 10^{-13}$
- Large transverse fluctuations from hadronic showers, long lived neutral emitted at large angles: similar to previous estimates
- Upstream interactions: requires precise knowledge of dead material in the beam line
  - SRD, V2, tracker suppression of secondaries
  - HCAL: lateral E and time spread compared with that expected from single electrons interacting in the ECAL target
  - estimation from data control regions
- Particle in-flight decays
  - SRD, ECAL energy and incoming track angle

# Backgrounds

Background source	Estimated number of events, $n_b$
hermeticity: punchthrough $\gamma$ 's, cracks, ..	$< 0.001$
loss of hadrons from $e^- Z \rightarrow e^- + \text{hadrons}$	$< 0.001$
loss of muons from $e^- Z \rightarrow e^- Z \gamma; \gamma \rightarrow \mu^+ \mu^-$	$0.005 \pm 0.001$
$\mu \rightarrow e \nu \nu$ , $\pi$ , $K \rightarrow e \nu$ , $K_{e3}$ decays	$0.02 \pm 0.004$
$e^-$ interactions in the beam line materials	$0.09 \pm 0.03$
$\mu$ , $\pi$ , $K$ interactions in the target	$0.008 \pm 0.002$
accidental SR tag and $e^-$ from $\mu$ , $\pi$ , $K$ decays	$< 0.001$
Total $n_b$	$0.12 \pm 0.04$

- Dominant contribution from upstream interactions
- 30% uncertainty also mainly due to upstream interactions
- Estimated from extrapolation of background control regions to signal region

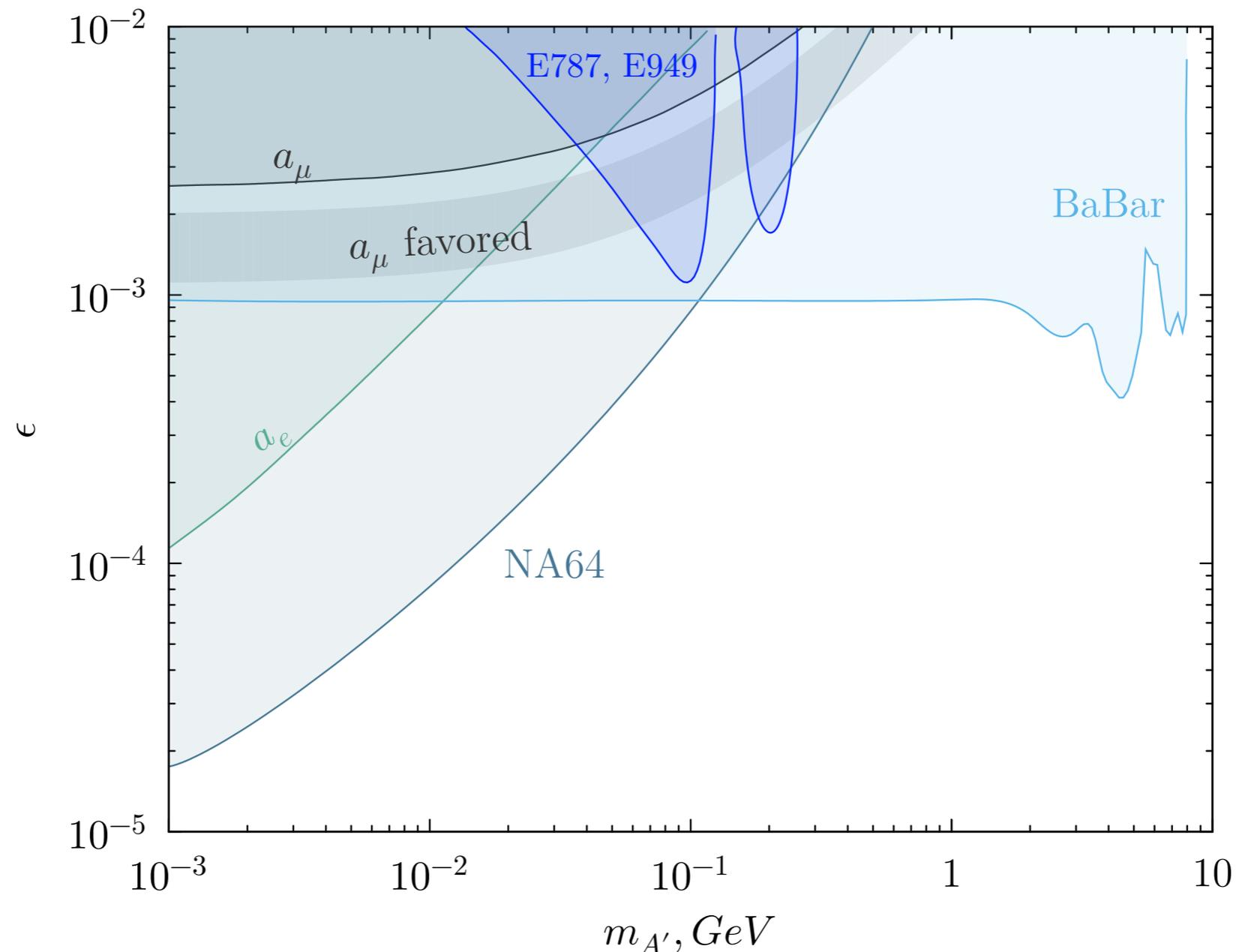
# Analysis

- Data collected from 2016 runs are divided in 3 bins: low, medium and high intensity beam.
- For each bin the background, efficiency corrections and uncertainties are estimated.
- A cut optimisation for the maximum sensitivity was performed for ECAL cut.
- The expected sensitivity was calculated with the Profile Likelihood method with RooStats, using the PL as test statistics, and taking the asymptotic approximation.

$$N_{A'} = \sum_{i=1}^3 N_{A'}^i = \sum_{i=1}^3 n_{EOT}^i \epsilon_{tot}^i n_{A'}^i(\epsilon, m_{A'}, \Delta E_e)$$

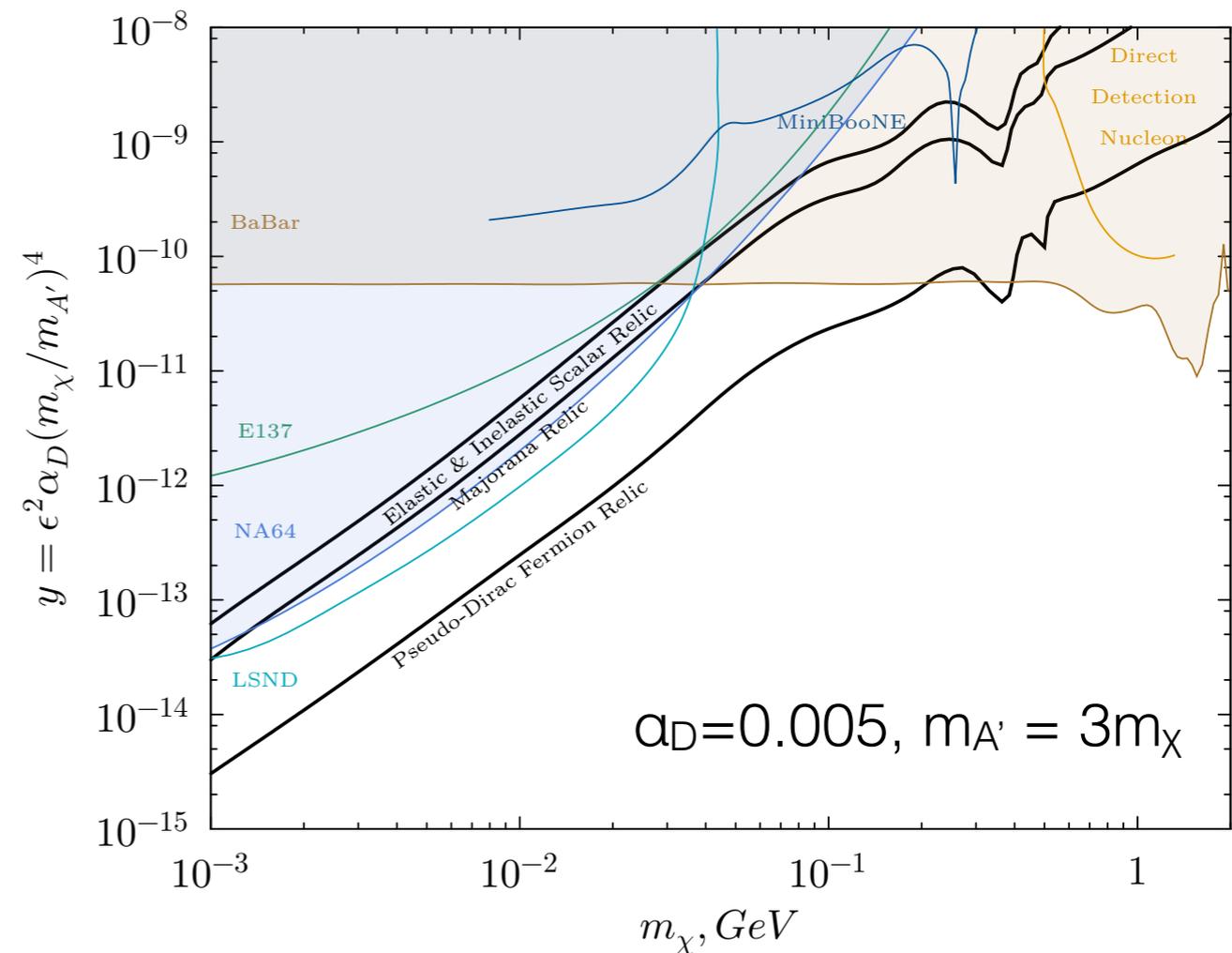
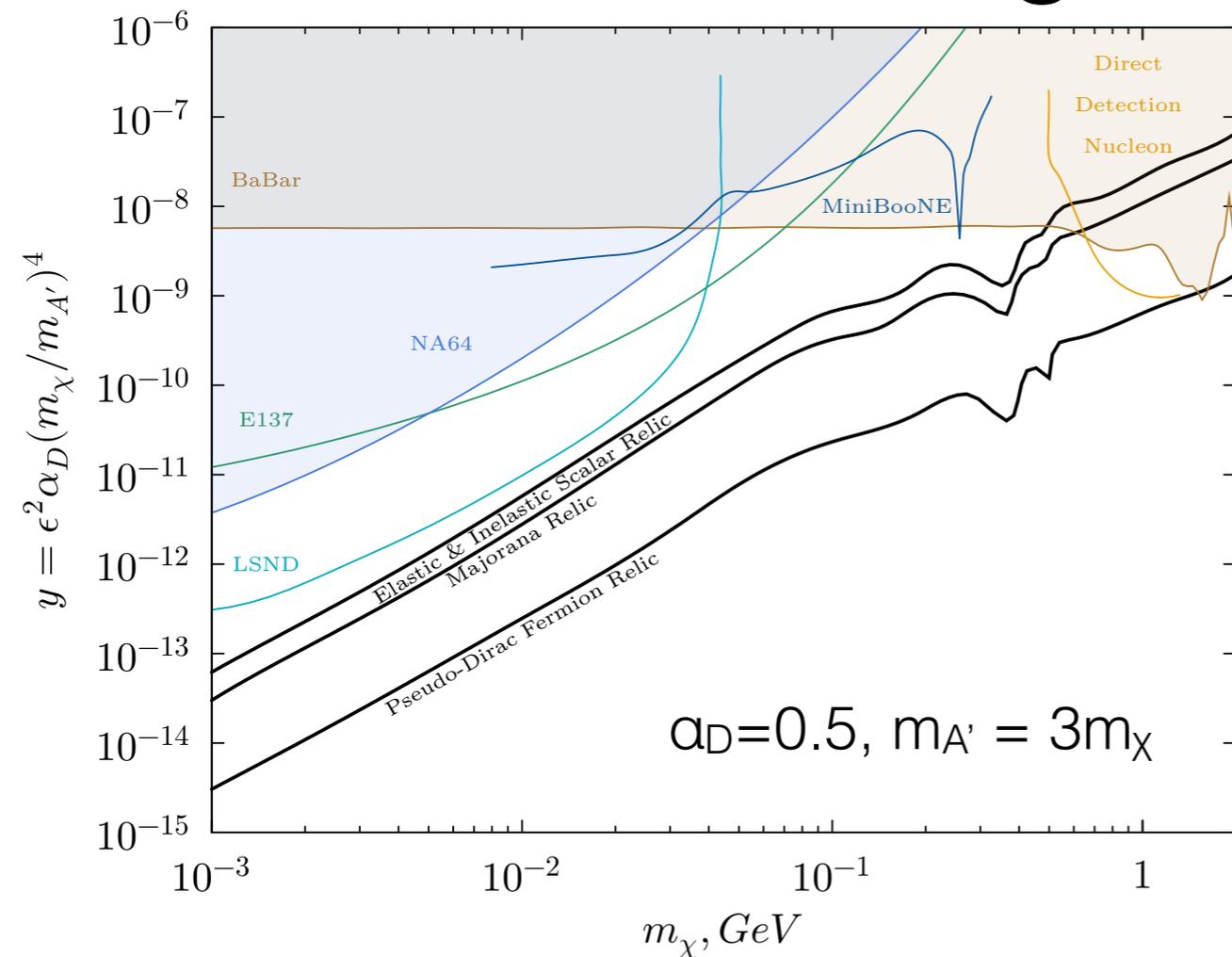
- Each  $i$ th entry for each data: simulating signal events for beam conditions and reconstructing w/ selection criteria, and efficiency corrections.
- Results also cross checked with simple limit from Poisson signal model with log-normal used for systematic uncertainty terms. Results agree within %.

# Results on $A'$ parameters



- Best limits in the region 0.001 - 0.1 GeV.
- Muon  $g-2$  favoured parameter region for vector mediator model excluded.
- Phys. Rev. Letters **118**, 011802 (2017)

# Results on light thermal dark matter



- LTDM models can be classified into spin and mass of DM and mediators, here only considering vector mediator.
- Assuming limits from prev. slide, constraints on DM annihilation freeze out.
- Results obtained for LSND, E137 and MiniBoone with  $10^{22}$ ,  $10^{19}$  and  $10^{20}$  POT.
- NA64 obtained with only  $\sim 4 \times 10^{10}$  EOT. With  $\sim 4 \times 10^{11}$  EOT NA64 can cover all beam dump exclusion areas.

# Conclusions

- Search is performed for sub-GeV dark photon mediated production of dark matter by NA64, using  $4.3 \times 10^{10}$  100 GeV electrons.
- No evidence of such events found.
- Derived upper limits on  $A'\text{-}\gamma$  mixing strength in the mass range 1-500 MeV, allowing to exclude vector mediator model solution for the muon  $g\text{-}2$  anomaly.
- Assuming these limits and constraints on DM ann. freeze out NA64 managed to exceed also limits on LTDM scenarios.
- NA64 continues to increase statistics in the near future and extend searches for dark matter and new physics at CERN SPS.
- Just finished our 2017 run, collecting additional  $5 \times 10^{10}$  electrons:
  - Runs finished both with invisible and visible mode, sensitivity to exclude  $\varepsilon = [5 \times 10^{-5}, 10^{-3}]$ , covering light X boson ( ${}^8\text{Be}$ ) favoured parameter region
  - Data under evaluation